C-terminal interactions mediate the quaternary dynamics of αB-crystallin


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αB-crystallin is a highly dynamic, polydisperse small heat-shock protein that can form oligomers ranging in mass from 200 to 800 kDa. Here we use a multifaceted mass spectrometry approach to assess the role of the C-terminal tail in the self-assembly of αB-crystallin. Titration experiments allow us to monitor the binding of peptides representing the C-terminus to the αB-crystallin core domain, and observe individual affinities to both monomeric and dimeric forms. Notably, we find that binding the second peptide equivalent to the core domain dimer is considerably more difficult than the first, suggesting a role of the C-terminus in regulating assembly. This finding motivates us to examine the effect of point mutations in the C-terminus in the full-length protein, by quantifying the changes in oligomeric distribution and corresponding subunit exchange rates. Our results combine to demonstrate that alterations in the C-terminal tail have a significant impact on the thermodynamics and kinetics of αB-crystallin. Remarkably, we find that there is energy compensation between the inter- and intra-dimer interfaces: when one interaction is weakened, the other is strengthened. This allosteric communication between binding sites on αB-crystallin is likely important for its role in binding target proteins.

1. Introduction

αB-crystallin is an oligomeric vertebrate small heat-shock protein (sHSP) with molecular chaperone activity [1]. In line with the general mechanism of action suggested for the sHSPs [2–5], αB-crystallin can trap non-native proteins in such a manner that prevents their aggregation [6] and facilitates their recovery by the downstream ATP-dependent molecular chaperone machinery [7]. Furthermore, αB-crystallin has been shown to bind to mature amyloid fibrils, inhibiting their elongation [8–11]. These functions represent an important mechanism that enables the cell to cope with the burden of unfolded proteins and maintain ‘proteostasis’ [12–14]. It is therefore no surprise that malfunction of αB-crystallin has been linked to numerous protein deposition diseases ranging from cataract formation [15] to motor neuropathies [16] and neurodegenerative disease [17].

Despite several thousand sHSP genes having been deposited in the UniProt database, there is relatively sparse detailed information about their assembled structures [2–5]. This is primarily due to their dynamic and frequently polydisperse nature, properties that make high-resolution studies particularly challenging [18]. αB-crystallin assembles into an ensemble of interconverting oligomeric states spanning approximately 10–40 subunits [19,20]. In recent years, considerable strides have been made in overcoming this heterogeneity and interrogating the structure of αB-crystallin. X-ray crystallography studies of N- and C-terminally truncated constructs of both αB-crystallin and its eye-lens-specific isoform αA [21–24] as well as solid-state nuclear magnetic resonance spectroscopy (ssNMR) of the full-length protein [25] have revealed the structure of a dimeric ‘building block’. This protomer is composed of a β-sandwich ‘α-crystallin’ core that assembles through anti-parallel (AP) pairwise interactions between extended β6 + 7 strands (figure 1). This interface
The crystal structure of a construct of αB-crystallin, truncated of the extension in addition to the N-terminal region, formed a runaway domain-swapped polymer, with the IXI lodged in the β4–β8 groove (figure 1b) [23]. This binding of the IXI was also demonstrated to pertain to the full-length protein in ssNMR performed at low temperature [25]. By contrast, NMR experiments performed in solution at physiological temperatures and pH revealed the IXI in full-length αB-crystallin to primarily populate an intrinsically disordered conformation [31,35–37]. The apparent contradiction between these studies has recently been rationalized by a temperature-dependent transition between conformations: a mixture of both bound and unbound states were observed below 0°C (in both solution and ssNMR experiments), whereas at temperatures above this only the unbound state was readily observed [36].

Fluctuations in the C-terminus on the millisecond timescale have been shown to be rate-limiting in the movement of subunits between αB-crystallin oligomers [31]. Moreover, removal of residues from the flexible C-terminal extension has been shown to reduce the rate of this subunit exchange [38,39]. Considering the tight regulation of αB-crystallin dynamics by residues in this region of the protein, and the seemingly crucial role of plasticity in sHSP function [2–5], it is perhaps not surprising that mutations near the C-terminus seemingly crucial role of plasticity in sHSP function [2-5], it is perhaps not surprising that mutations near the C-terminus display aberrant molecular chaperone function in vitro [32,40], in cells [41,42] and are associated with disease [43–46].

In this study, we investigate the molecular details of the interaction between the α-crystallin C-terminal region and the β4–β8 groove in the core domain. We perform our investigations using nano-electrospray mass spectrometry (MS) and ion-mobility spectrometry (IM). Since its initial applications to the study of protein assemblies in the early 1990s [47], MS has matured as a structural biology approach to allow the accurate determination of oligomeric stoichiometry and dynamics [48–50]. We exploit these benefits here to characterize in detail the influence of residues in this region of the protein, and the dynamics by residues in this region of the protein, and the dynamics by residues in this region of the protein.
forming the more tightly bound conformations observed in other sHSPs. To extend this work to the context of the full-length protein, we describe a series of point mutations to alanine (figure 1c), and characterize their effect on the oligomeric distribution and subunit exchange dynamics of the protein. We find that mutations in the C-terminal tail significantly affect the stability and dynamics of the oligomers, indicating that, while predominantly disorderd, they are nonetheless able to exert significant influence on determining which oligomers are populated. We note that mutations of pseudo-equivalent palindrome residues do not display equivalent changes in thermodynamics and kinetics, suggesting a favoured orientation within the oligomers. Moreover, we observe a strong negative correlation between the free energy of the inter- and intra-dimer interface, revealing an allosteric coupling between them.

2. Materials and methods

(a) Protein expression and purification

Full-length αβ-crystallin, the C-terminal mutants and a truncated form (residues 68–153, αβ_{Trunc}) were expressed in *Escherichia coli* and purified as previously described [23,51]. The peptide, ERTIPITRE, was expressed and purified from *E. coli* as described previously [52], whereas peptides ERTIPITREETPAVK and TIERPREIT were purchased from BioMatik (Canada). Unless otherwise stated, solutions of full-length αβ-crystallin were prepared at a monomeric concentration of 20 μM in 200 mM ammonium acetate pH 6.9 prior to MS analysis.

(b) Measuring peptide binding to the αβ-crystallin core domain

All titration experiments were carried out at room temperature and allowed to equilibrate for 30 min prior to measurement. Solutions of αβ_{Trunc} were prepared at a final monomeric concentration of 4 μM (as determined using a bicinchoninic acid assay on a stock solution) in 200 mM ammonium acetate pH 6.9, whereas the final concentrations of the C-terminal peptides were varied. All C-Terminal peptides were dissolved in water, with final concentrations of 4, 8, 16, 32, 64, 128 μM.

Titrations were collected on a Synapt G1 HDMS instrument (Waters Ltd., UK), modified to incorporate a linear drift tube [53]. Parameters were set to minimize activation of gas-phase ions and maximize separation between monomeric and dimeric species in the arrival-time dimension. To ensure the preservation of the labile interaction between peptide and αβ_{Trunc} in the mass spectrometer, and reduce the number of alkaline metal adducts, a small reservoir of acetonitrile was maintained in the source region of the mass spectrometer [54]. The instrumental voltages used were capillary (1400), sampling cone (10), extraction cone (1), trap (5), bias (20) and drift tube (50). The gas flow rates were trap (argon, 2.6 ml min⁻¹), drift tube (nitrogen, 50 or 20 ml min⁻¹), and the ‘backing’ pressure was 3.8 mbar. Three spectra were recorded for each titration point, using a new needle each time to account for needle-to-needle variation.

(c) Ion-mobility mass spectrometry titration data analysis

External calibration of the spectra was performed using MassLynx software (Waters Ltd.), and the data were exported for analysis using our home-built algorithm CHAMP [55]. CHAMP was adjusted to incorporate the arrival-time dimension, along the lines we have described previously [33]. The collisional cross-section (CCS) of each molecular species was considered to be invariant with charge state, an assumption that appears valid for protein assemblies in the absence of peturbants [56].

To ensure the stability of the protein assemblies in the gas phase, and CCSs close to those expected from their solution structure, we examined them in a charge-reduced form [57], achieved here by a partial pressure of acetonitrile in the source region [54]. Individual IM–MS peaks were considered to have Gaussian profiles in the arrival-time dimension [33], and asymmetric Lorentzian profiles in the mass-to-charge (m/z) dimension [58]. The width of the IM peaks was fitted by assuming the full-width-at-half-maximum scales linearly with drift time for ions of similar mobility [59]. Minimization was performed using a combination of stimulated annealing and Levenberg–Marquardt methods [55].

(d) Determination of peptide binding affinity

To determine the binding affinities of the C-terminal peptides to αβ_{Trunc}, we used the abundances extracted by CHAMP as described earlier. In order to obtain reliable binding affinities, it is important to account for ‘false-positive’ complexes [60]. These arise from non-specific association during the electrospray process, and can become significant at high protein or ligand concentrations [61]. To address this, we performed two sets of control experiments to assess the extent of non-specific binding. In the first, we obtained spectra of the ‘scrambled’ C-terminal peptide TIERPREIT incubated at a range of concentrations with αβ_{Trunc}, and in the second, the specific peptide ERTIPITRE with the unrelated protein cytochrome c. In both cases, the abundance of protein–ligand complexes can be entirely ascribed to non-specific association. In this way, the abundances of detected protein–ligand complexes can be corrected to reflect solely the specific associations [62]. Both these control experiments were found to correspond to an apparent K_d on the order of millimolar, see the electronic supplementary material. Thus adjusted, the experimental distribution of apo- and holo-αβ_{Trunc} was examined in terms of a protein–ligand binding model, described in detail in the electronic supplementary material, in order to obtain the binding affinities.

(e) Assessment of residue-specific C-terminal binding in silico

A model structure for ΔG calculations was constructed from the truncated human αβ-crystallin crystal structure (PDB: 3L1G) with the C-terminal tail bound [23] and further modified using atomic coordinates from the bound tail in the structure of truncated zebrafish αα-crystallin (PDB: 3N3E) [24]. Atomic coordinates from 3L1G for residues 68–153 and residues 153–162 of the bound C-terminal tail from a symmetry-derived molecule were merged into one chain. The C-terminal tail was modified by grafting residues 163–166 of 3N3E after alignment of the tail. This model was relaxed using *R*seta v. 3.3, prior to estimating ΔG_s upon mutation in *silico* to alanine [63,64].

(f) Determination of the oligomeric distributions and interface free energies

All samples were pre-incubated at 37 °C for 30 min prior to mixing to ensure they had reached equilibrium. MS was performed under activating conditions as described previously [65]. The resulting spectra can be unambiguously ascribed to the oligomeric distribution in solution owing to the predictable dissociation pathway of protein assemblies in the gas phase [66]. The relative abundances of all the doubly stripped...
oligomers were quantified from the intensity of the peaks [19], and fitted to an oligomerization model we have described in detail elsewhere [65] to obtain accurate measurements of the free energy of the inter- and intra-dimer interfaces, $\Delta G_s$ and $\Delta G_d$ respectively ($\Delta G_{s+d} = \Delta G_s + \Delta G_d$) (see the electronic supplementary material, table S1). In order to compare the relative effects of the mutations in their various positions, the free energies were considered as $\Delta \Delta G_s$ according to $\Delta \Delta G = \Delta G_{\text{mutant}} - \Delta G_{\text{WT}}$.

(g) Monitoring quaternary dynamics of $\alpha B$-crystallin
Subunit exchange experiments were performed by incubating wild-type and mutant forms of $\alpha B$-crystallin with the $^{13}$C isotopically labelled $\alpha B_{ERT}$, with the exception of WT protein that was mixed with its labelled equivalent. All samples were pre-incubated individually at 37°C for 30 min to ensure they had reached an equilibrium oligomeric distribution prior to mixing at a 1:1 ratio. Aliquots of the mixture were taken at various time points, with the reaction quenched on ice, and analysed off-line as described previously [65]. The peak corresponding to all oligomers, each carrying the same number of charges as subunits (i.e. $\alpha B$, ); figure 2), enable the time-dependent disappearance of homo-oligomers, and concomitant emergence of hetero-oligomers, to be readily monitored [38,67]. The data were then fitted to our oligomerization model [65], allowing the rate constants of dissociation $k_-$ and $k_-$, and rate of association $k_+$ to be extracted (see the electronic supplementary material, table S1).

3. Results

(a) The C-terminal tail of $\alpha B$-crystallin bind the core domain only weakly
To examine the binding of the tail residues of $\alpha B$-crystallin, we investigated the binding of peptides to a truncated construct $\alpha B_{Dom}$, residues 68–153, in which the $\beta 4–\beta 8$ groove is necessarily unoccupied. We performed titration experiments in which we incubated $\alpha B_{Dom}$ with either the palindromic peptide ERTIPITRE ($\alpha B_{Pal}$), representing residues 156–164 of $\alpha B$-crystallin, or the peptide ERTIPITREEPKAVTAAPKK ($\alpha B_{Ext}$), residues 156–175, the combined C-terminal tail and extension (figure 1d).

At the lowest concentration of peptide (4 $\mu$M), peaks are observed corresponding to $\alpha B_{Dom}$ monomer and dimer (figure 2a). The preponderance of monomer is in line with a weak intra-dimer interface [68] and $K_1$ in the low micromolar range [23]. Upon increasing the concentration of the peptide, peaks appear that correspond to the complexes formed between the protein and peptide at stoichiometries of ($\alpha B_{Dom}$)$_1$($\alpha B_{Pal}$)$_1$, ($\alpha B_{Dom}$)$_2$($\alpha B_{Pal}$)$_1$, and ($\alpha B_{Dom}$)$_2$($\alpha B_{Pal}$)$_2$ (figure 2a). Notably, even at the highest peptide concentrations for which we were able to obtain good mass spectra (128 $\mu$M peptide, more than a 10-fold excess), we never observed more than about half of the available sites to be occupied. Analogous experiments with $\alpha B_{Ext}$ showed apparently similar levels of bound forms (figure 2b). Binding of the C-terminal peptide therefore appears to be relatively weak, in agreement with previous data [31,36,37,69].

(b) Ion-mobility spectrometry—mass spectrometry allows quantitative extraction of peptide binding affinities to the $\alpha B$-crystallin core domain
While these MS data show qualitatively that the peptides bind to $\alpha B_{Dom}$, they are challenging to interpret quantitatively owing to the overlap of charge states from the different species, for example ($\alpha B_{Dom}$)$_4$$^+$ and ($\alpha B_{Dom}$)$_8$$^+$. This is particularly noticeable in the case of the cluster of peaks around 2470 m/z for $\alpha B_{Dom}$ incubated with $\alpha B_{Ext}$ (shaded, figure 2b). In order to facilitate the deconvolution of these peaks, we used a hybrid IM–MS strategy. These experiments provide an orthogonal dimension of separation to m/z according to the ions’ ability to traverse a tube of neutral gas at low pressure under the influence of a weak electric field. The transit time through the drift cell of the ions is proportional to their CCS and inversely proportional to charge [70].

An IM–MS spectrum of $\alpha B_{Dom}$ incubated with $\alpha B_{Ext}$ shows a series of charge states separated in both m/z and arrival time (figure 3a). The separation between monomer and dimer is dramatically improved relative to the one-dimensional MS experiment (figure 2). In order to analyse...
In this spectrum, we used an extension of our spectrum calculation algorithm CHAMP [55], modified to accommodate the arrival-time dimension [33]. By calculating different IM–MS spectra from candidate distributions of the molecular components in the spectrum, and comparing them to the data, a best-fit spectrum is obtained (see figure 3b and electronic supplementary material, figure S1). The correspondence between data and fit is excellent, allowing us to deconvolve the relative contributions of the different stoichiometries to the overall spectrum (figure 3c).

Using this approach, we analysed our titration series in order to interrogate the concentration dependence of peptide binding (see the electronic supplementary material, figure S1). At the lowest concentration of peptide (4 μM), peaks are visible corresponding to αB<sub>a</sub>-x monomers and dimers, and negligible binding of peptide is observed for either αB<sub>a</sub>-x or αB<sub>a</sub> (figure 4a, upper panels). At the highest concentration (128 μM) however, significant abundances of the bound forms are present (figure 4a, lower panels). It is notable how in the case of αB<sub>a</sub> significantly less (αB<sub>a</sub>)<sub>2</sub>(Pep)<sub>2</sub> is observed than for αB<sub>a</sub>-x.

To interpret these differences, we examined the abundances of the different molecular species in the context of a protein–ligand binding model (see electronic supplementary material, figure S1). For both αB<sub>a</sub>-x and αB<sub>a</sub>, we find that the K<sub>D</sub> of binding the first peptide equivalent was found to be approximately 70 μM, with a dimer interface K<sub>D</sub> of approximately 2 μM (orange and brown, respectively, figure 4b,c). Remarkably, the affinity for the second peptide equivalent we find to be substantially lower: 150 μM for αB<sub>a</sub>-x and 300 μM for αB<sub>a</sub> (grey, figure 4b,c). Furthermore, the dimer interface for (αB<sub>a</sub>)<sub>2</sub>(Pep)<sub>2</sub> is slightly weaker (4 μM for αB<sub>a</sub>-x and 6 μM for αB<sub>a</sub>) than in the other dimeric species. Taken together, our data demonstrate that binding two peptides to a dimer is substantially more difficult than binding just one, and has the concomitant effect of actively weakening the intra-dimer interface. This destabilization effect of the second peptide is substantially greater for αB<sub>a</sub> relative to αB<sub>a</sub>-x. Extrapolated to the context of oligomers of full-length protein therefore, the concurrent binding of two C-termini would precipitate the dissociation of a monomer, in a manner consistent with our previous NMR measurements [31,35,36].

(c) *A mass spectrometry assay to determine the thermodynamic and kinetic consequences of point-mutation*

To further analyse the interactions between the C-terminus and the α-crystallin domain in the context of the oligomer, we performed an alanine scan, generating a series of point mutants of full-length αB-crystallin that collectively act to encompass the entire palindromic region of sequence ERTIPITRE<sup>164</sup> (figure 1), and examined them by means of MS. Our method has been detailed previously [65], and is described briefly below. Using appropriate ion generation and transmission conditions [71], the intact αB-crystallin oligomers are directly measured in the mass spectrometer. The resultant mass spectra feature a broad region of signal arising from a large number of stoichiometries with many overlapping charge states [19], that is essentially uninterpretable (figure 5a). To overcome this, we use a collisional activation strategy in which the oligomer ions are heated by successive collisions with a target gas until they dissociate [72]. The general mechanism of gas-phase dissociation dictates that oligomers dissociate into highly charged monomers (figure 5b, 1000–2000 m/z) and complementary ‘stripped oligomers’ (figure 5b, 16 000–24 000 m/z) [66]. From these data, we can assign the different peaks, extract their intensities, and thereby determine the oligomeric distribution of the αB-crystallin ensemble (figure 5c). This approach is cross-validated by the faithful back-calculation of the original mass spectrum [33,55]. Further verification comes from noting that, while of much higher resolution, this distribution matches that measured by using solution phase methods extremely well [48].

Although αB-crystallin populates many stoichiometries, their distribution can be well described using a relatively...
simple oligomerization model [65]. This model invokes only two interactions between individual αB-crystallin monomers, corresponding to intra-dimer (dimer, d) and inter-dimer (edge, e) interfaces (see schematic in the electronic supplementary material, table S1). This two-parameter model assumes individual oligomers are in dynamic equilibrium with their corresponding monomers, an assumption justified by the appearance of oligomers comprising an odd number of subunits in experimental data and in facile subunit exchange [65]. The model predicts that the basic monomeric structure is identical in all oligomers, which is confirmed by NMR experiments [25,31], and that the dimer interface is labile, as suggested by our results on the core domain (figure 2) and previous measurements [23,68,73]. Through fitting this model to the MS data, the oligomeric distributions can be reproduced accurately, and the association free energies of the edge and dimer interfaces (ΔG_{e}, ΔG_{d}, respectively) quantified (figure 2d) [65].

To complement these thermodynamic parameters, the rate constants that govern the inter-conversion of the αB-crystallin oligomers can be determined using subunit exchange reactions in which the protein is incubated with a labelled counterpart [74]. This is illustrated schematically for the reaction between αB-crystallin and its 13C-equivalent (figure 5e). The peak at ≈20 100 m/z (figure 5e, black) comprises αB-crystallin-striped oligomers, each carrying as many charges as subunits, and is representative of all oligomers that contribute to the ensemble [19]. Mass spectra of the 13C protein reveal the same oligomeric distribution as the 12C counterparts, but, due to the additional mass, the equivalent peak appears at ≈21 000 m/z (figure 5e, purple). Upon incubation, the two peaks corresponding to homo-oligomers coalesce into one broad peak centred on their midpoint (figure 5f, green), indicating the formation of hetero-oligomers [38]. This subunit exchange can be monitored by collecting data at different time points (figure 5f, red), and fitting these data to simulated time courses generated using our oligomerization model (figure 5f, blue) allows the extraction of the corresponding rate constants [65].

(d) Destabilization of the C-terminal interaction stabilizes the intra-dimer interfaces

We obtained the oligomer distributions for the wild-type protein and the seven different mutants (nomenclature given in figure 1) at pH 7 and 37 °C. In all cases, the proteins assemble into polydisperse ensembles centred around a 24–28 mer (figure 6a), consistent with previous measurements for αB_{WT} and phosphorylated versions [65,75]. In addition, all of the distributions were found to have a ‘saw-toothed’

![Figure 4](http://rstb.royalsocietypublishing.org/Downloaded from http://rstb.royalsocietypublishing.org/ April 19, 2017)

Figure 4. (a) Representative IM–MS spectra of αB_{Dom} incubated with either αB_{Pal} (left panels) or αB_{Ext} (right panels). At 4 μM peptide (upper panels), both αB_{Dom} monomers and dimers are observed, with very little peptide binding, revealing that the K_D must be significantly larger than 1 μM. At 128 μM peptide (lower panels), both apo and holo forms are observed. Notably, the abundance of (αB_{Dom})(Pep)_{2}, the domain dimer with two peptides, is less abundant in the case of αB_{Dom} than αB_{Pal}, whereas the quantities of (αB_{Dom})(Pep), are similar in both. (b) Quantitative analysis of peptide binding. We analyse the MS titration data in terms of a ligand binding model invoking a total of six coupled equilibria (see the electronic supplementary material for further details) with an increasing number of free parameters. In the simplest model all sets of monomer/dimer equilibria (brown, grey) and peptide binding equilibria (orange, blue) are identical so that the system is controlled by only two independent parameters. Allowing the second peptide binding to have different binding affinity to the first binding event leads to a model specified by four affinities (orange, blue, brown and grey), determined from three independent fitting parameters. The inclusion of this additional parameter is statistically justified for the titrations with αB_{Pal} and αB_{Ext}, whereas models of increased complexity were found not to be. As described in the text, these data enable us to conclude that the second peptide binds with lower affinity than the first. The raw data and fits are shown in the electronic supplementary material, figure S1. (c) A bootstrap analysis was performed to assess the uncertainty in our estimates of dissociation constants as described in detail in the electronic supplementary material. The resulting histograms are shown for fitting the statistically justified four K_D values to data from αB_{Pal} (solid lines, colours as in b) and αB_{Ext} (dashed lines, colours as in b). The histograms reveal that a single peptide has a similar affinity to αB_{Dom} in both cases (orange). The affinity of binding the second peptide to αB_{Dom} is lower for αB_{Ext} than αB_{Pal} (grey). Moreover, the K_D of (αB_{Dom})(Pep)_{2} is notably higher in the case of αB_{Ext} than αB_{Pal}, indicating weaker binding.
character, with oligomers comprising an even number of subunits more abundant than those composed of an odd number. Significantly, the extent of this disparity varies according to the location of the mutation: for example, the αBnrt mutant is ‘spikier’ than αBnvt. In the context of the model, increased disparity is diagnostic of a strengthening of the dimer interface, and vice versa.

In order to quantify this effect, we determined \( \Delta G_a \) and \( \Delta G_d \) values for each of the proteins (see the electronic supplementary material, table S1). In order to readily compare the relative effects of the mutations, we calculated the corresponding \( \Delta \Delta G_s \), where a positive value suggests that the mutant is less stable than the wild-type, and a negative value the converse (see the electronic supplementary material, table S1). A plot of \( \Delta \Delta G_d \) versus \( \Delta \Delta G_s \) reveals a clear negative correlation of these two quantities (figure 6b). The majority of the mutations cause a weakening of the edge interface, of up to \( 2 \ \text{kJ mol}^{-1} \). Remarkably, we observe a concomitant and compensatory strengthening of the dimer interfaces. This demonstrates a clear allosteric coupling between the two interfaces [76]. Tighter binding of the C-terminal tail to the β4+8 groove weakens the interaction across the intra-dimer interface formed by the β6–7 strands, and vice versa. This finding is consistent with our titration data on αBnvt, where we show that binding the second petide to the dimer destabilizes the dimer interface. Furthermore, our results reveal that residues upstream of I161 have significant effects on the distribution, whereas those downstream of I161 (as exemplified by αBrr) have negligible consequences on the distribution.

Structures of isolated αB-crystallin dimers have been determined in which the C-terminus is bound into the β4–8 groove. It is possible to estimate the difference in free energy upon point mutation \textit{in silico}, allowing us to obtain theoretical \( \Delta \Delta G_s \) values to compare with our data (see the electronic supplementary material, table S1). While there is some qualitative correspondence in that most of the mutations are predicted to destabilize the interaction both \textit{in silico} and \textit{in vitro}, quantitatively there are dramatic differences in the magnitude of the values and discrepancies as to which residues have the greatest contribution. Most notably, the \textit{in silico} calculations on the bound state predict that the mutations to the IXI region should cause a destabilization of almost \( 40 \ \text{kJ mol}^{-1} \), an effect approximately 30-fold greater than that measured experimentally. This apparent conflict might however be rationalized by the IXI populating a bound state to the order
of only a few per cent, consistent with solution-state NMR experiments [31,38].

Furthermore, the calculations predict the αB_{TRE} mutation to cause a significant destabilization of the interaction. By contrast, we see no change in the free energy, within the limits of experimental error, for this triple-mutant. Taken together, we can conclude that the free energies predicted for the bound state of the C-terminal tail are very far removed from that measured in solution. By contrast, the calculations are broadly consistent with our, and others’ [69], titration experiments on αB_{Dom} that reveal K_D in the 70–300 μM range; and our NMR measurements that show that in the oligomers the tail is largely present in a disordered, and unbound conformation [31,38].
(e) C-terminal mutations cause an increase in both
dissociation and association rates
In order to complement our measurement of the thermodynamic consequences of the mutations, subunit exchange experiments were performed to obtain the monomer association and dissociation rates. Incubations between unlabelled and 13C protein were monitored as described above, and the time taken for the respective homo-oligomers to equilibrate into a distribution of hetero-oligomers was determined. For αBcryst, this process was complete after approximately 40 min (figure 6c). By contrast, all the mutants exchanged significantly faster, with the fastest, αBcryst-

BDom, reaching equilibrium within 20 min (figure 6c). This demonstrates that mutations in the C-terminal region affect not only the strength of the quaternary interfaces within αB-crystallin, but also their associated dynamics.

The difference in subunit exchange rates was quantified by extracting the dissociation rate constants $k_\text{a}$ and $k_\text{d}$ and pseudo-first-order association rate $k_\text{a}$, where $[\text{B}]_1$ is the concentration of free monomer in solution (see the electronic supplementary material, table S1). In the case of the mutations studied here, all display both increased association and dissociation rates. To enable comparison between the mutant and the wild-type, we calculated $\Delta \Delta G^\ddagger$ values (change in the free energy of activation) in each case. In all cases, the $\Delta \Delta G^\ddagger$ are significantly negative and, similar to the finding with the thermodynamic data (figure 6b), $\Delta \Delta G^\ddagger$ (association) and $\Delta \Delta G^\ddagger_{a,d}$ are also strongly correlated (figure 6d). In other words, all mutations that increase the rate of association increase the rate of dissociation by approximately the same amount. Similar to the result obtained for the thermodynamic parameters, mutations upstream of I161 were found to have the most significant effects.

4. Discussion
We have presented a detailed investigation into the influence of the palindromic C-terminal residues on the thermodynamics and kinetics of αB-crystallin oligomers. We have addressed the problem using two orthogonal approaches. First, we performed titration experiments between peptides mimicking the C-terminal region of αB-crystallin and αBcryst-

BDom, determined their binding affinities, and noted effects on the intra-dimer interface induced by binding. Second, we mutated residues in the C-terminal tail and observed the thermodynamic and kinetic consequences on the oligomers. Both approaches give results that are internally consistent and enable a deeper understanding of the molecular interactions that dictate the properties of αB-crystallin oligomers.

(a) Mass spectrometry for studying the
thermodynamics and kinetics of protein interfaces
We employed an IM–MS strategy to analyse titrations of αBcryst-

BDom and C-terminal mimicking peptides to robustly quantify the different molecular species that coexist in solution. This allowed us to extract the $K_{\text{d}}$ of peptide binding and, rather than providing an ensemble average over the stoichiometries present at equilibrium, do so for both the monomeric and dimeric forms of αBcryst-

BDom. While quantitatively consistent with ensemble measurements from NMR [69], interestingly we observe that two sequential peptide binding events to (αBcryst-

BDom)2 are not equivalent. While binding of the first peptide is itself weak, it is not only significantly more difficult to bind the second but there is also a concomitant destabilization of the dimer interface. This effect is particularly pronounced in the case of αBcryst-

BDom, which is longer than αBcryst-

BDom and has a significant number of positively charged residues at the C-terminus. Both steric and coulombic interactions are therefore likely to play a role in decreasing the binding affinity of the second peptide. As such, while one might intuitively expect the high local concentration of C-terminal tail in the oligomer to effectively overcome its weak binding affinity [69], our results indicate that additional tail-binding events are less favourable and precipitate dissociation of the oligomer. This is supported by the observation that not only does truncation of the C-terminus reduce the rate of subunit exchange of the α-crystallins [38,39], but also that mutating charged residues affects their self-assembly [40]. Binding of the second peptide has an additional interesting effect in that the intra-dimer interface is destabilized. This remarkable finding is consistent with the model derived by solution-state NMR that suggests that subunit exchange is facilitated by the C-terminus binding the β4–8 groove [31,36].

To complement studies on the truncated αB-crystallin construct, we examined point mutants of the full-length protein. The effects of individual and multiple alanine mutations on both thermodynamic and kinetic properties of the oligomers were then quantified by means of MS. The advantages of MS for monitoring the quaternary consequences of such alanine-scanning rest in the high resolution of separation in both mass and time afforded by the approach [48]. As a result, in this study, we have successfully extracted the changes in rate constants and free energies using experiments performed on the minute time-scale.

(b) The inter- and intra-dimer interfaces of
αB-crystallin are allosterically coupled and
energy-compensate to maintain oligomer size
Our study has revealed that while the average oligomer size was essentially unchanged by mutation, the individual interactions that define the assembly are significantly impacted. We find that destabilization of the edge interface caused by alanine mutation results in a stabilization of the intra-dimer interface. This finding shows that the effect of binding the C-terminus at a first site (in this case, the β4–8 groove) influences the interactions at a second, distal site (the β6–7 dimer interface): the hallmark of allosteric communication within the protein [76]. This phenomenon was previously noted in experiments examining the influence of pH on the distribution of wild-type αB-crystallin [65]. Similarly, phosphorylation of the N-terminus, which weakens the dimer interface [75,77], likely by binding a cleft on the inside of the oligomer [22], results in a strengthening of the edge interaction [65]. This role of the β4–8 groove and C-terminal tail in allosteric communication rationalizes their identification as important regulatory regions in the molecular chaperone function of αB-crystallin [78–81].

Taken together, while the inter- and intra-dimer free energies can vary substantially, the effect of this energy compensation is such that the quantity $\Delta G_{e+d}$ is kept...
approximately constant (figure 6b). This is the value that pre-
dominantly dictates the average oligomer size [65], revealing
that αB-crystallin oligomers have the ability to radically
alter their interface dynamics, as well as structure [23], yet
retain an essentially constant gross oligomeric distribution.
Functionally, this is likely to be an important property for
maintaining eye-lens transparency at high protein concen-
tration while avoiding crystallization [82]. In addition, the
molecular chaperone activity of sHSPs may itself benefit
from such polydispersity, through the provision of a diversity
of binding surfaces for intercepting a wide range of destabi-
lized target proteins [83]. With polydispersity potentially
crucial to both these roles of αB-crystallin in the body, it is
perhaps unsurprising that the vast number of mutations
reported in the literature have had only a limited effect on
the overall oligomerization of the protein, and the quest for
a homogeneous quaternary structure remains unfulfilled [18].

(c) The synergistic roles of the C-terminal tail, extension
and palindrome

Our experiments have revealed the significant contributions
made to the oligomeric dynamics and interface stabilities
by residues in the C-terminal region of αB-crystallin. It is
notable that none of the mutations resulted in complete dis-
assembly of the oligomer, which, together with evidence
that truncated constructs of the α-crystallins remain assembly
competent [23], reveals that the C-terminus is not the sole
provider of thermodynamic stability to the oligomer. Instead,
its role seems to be more subtle, acting as a ‘gate-keeper’ for
the quaternary dynamics of the protein [31].

We have demonstrated here that all of the mutations on
the C-terminal tail have small but measurable effects
(ΔΔG < 2 kJ mol⁻¹) on the strength of the edge interface
and the corresponding rates of association and dissociation.
The effect of these changes is to modulate the proportion of
time during which the C-terminal IXI remains bound to the
oligomer. This provides a biophysical rationale as to why
mutation or removal of C-terminal residues leads to a faster
function of αB-crystallin in vitro [32,40], in cells [41,42],
and have been identified in dilated cardiomyopathy [44,45] and
cataract [43,46]. Furthermore, a recent report on the mono-
disperse archaeal HSP14.0 apparently indicated a similarly
transient C-terminal interaction [30], which raises the possi-
bility that a regulatory role for this region of sequence may
be widespread in the sHSP family.

While mutations upstream of I161 were found to signifi-
cantly destabilize the inter-dimer interface, mutations
downstream had no discernable effects on the distribution
(figure 7). Truncated α-crystallin constructs have been crystal-
lized with the C-terminus in two different orientations,
facilitated by its palindromic nature [23,24]. Our data show
that while mutation of the ERT before the IXI has signifi-
cant effects on the distribution, mutation of the TRE following the
IXI does not, suggesting that in solution, the oligomers prob-
able have a significantly preferred orientation. This is in line
with previous NMR results that have demonstrated the exten-
sion to tumble freely in solution [84], even in homogenates
mimicking the crowded environment of the eye lens [85] or
solid-state NMR preparations [36]. Furthermore, we have
found that the Kd of the αBext is significantly higher than
that of αBpref, demonstrating that the extension acts to
facilitate the detachment of the tail. This acts to directly
influence the rate of subunit exchange, demonstrating why
mutant proteins with shorter extensions have reduced quaternary dynamics [38,39].

5. Conclusions

Our study demonstrates that MS can be used to provide a
highly detailed view of the macroscopic quaternary structure
and dynamics of αB-crystallin. Moreover, they strongly sup-
port a model we recently proposed to explain many
thermodynamic and kinetic properties of the oligomers. We
show in both truncated αB-crystallin constructs and full-
length oligomers that there is allosteryc communication
between the β4–8 groove and the β6 + 7 dimer interface.
When two C-terminal peptides bind, the dimer interface is
significantly destabilized, thus facilitating monomer disso-
ciation and subunit exchange, and potentially exposing
target binding sites. Moreover, mutational studies reveal
that when tail-to-groove interactions are destabilized there
is a corresponding increase in strength in the dimer interface
within the oligomers and an increase in the rates at which
monomers associate and dissociate. These observations
provide a rationale as to why variant proteins with muta-
tions in this region are associated with multiple disease
states. Furthermore, our MS approach allows the quantifi-
cation of the quaternary dynamics and oligomerization of a
polydisperse protein, which not only opens the door for
assessing the consequences of disease-related point mutants
but also screening for small molecules that may act to prevent
or reverse this.

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